





Aspects of Graphene Nanoribbons Devices Simulations

Pedro Brandimarte

June 13, 2016









• High mobility \longrightarrow 10⁵ cm²V⁻¹s⁻¹



K. Novoselov et al. Science **306**, 666-669 (2004).









High mobility → 10⁵ cm²V⁻¹s⁻¹





• Gap absence:



A. Castro Neto et al. Rev. of Mod. Phys. 81, 109-162 (2009).







Graphene Nanoribbons (GNRs)

• Bottom-up fabrication of both AGNR and ZGNR by on-surface reaction of molecular precursors.



J. Cai et al. Nature **466**, 470 (2010).

L. Talirz, P. Ruffieux, and R. Fasel. Advanced Materials (2016).







Graphene Nanoribbons (GNRs)

 Bottom-up fabrication of both AGNR and ZGNR by on-surface reaction of molecular precursors.



J. Cai et al. Nature **466**, 470 (2010).

L. Talirz, P. Ruffieux, and R. Fasel. Advanced Materials (2016).

 Semiconductor character, with energy gap depending on their width and shape.



Y.-W. Son et al. Phys. Rev. Lett. **97**, 216803 (2006). Y.-C. Chen et al. ACS Nano **7**, 6123 (2013). A. Kimouche et al. Nature Communications **6**, 10177 (2015).





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Electron Transport Simulations of 4-Terminal Crossed Graphene Nanoribbons Devices

Pedro Brandimarte, Nick R. Papior, Mads Engelund, Aran Garcia-Lekue, Thomas Frederiksen, and Daniel Sánchez-Portal

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M. Masum Habib and Roger K. Lake. Phys. Rev. B 86, 045418 (2012).

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Crossed 14-AGNR





+

Non-Equilibrium Green's Function (NEGF)

TranSIESTA

J. M. Soler et al. J. Phys. Condens. Matter. **14**, 2745 (2002). Mads Brandbyge et al. Phys. Rev. B **65**, 165401 (2002).

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Crossed 14-AGNR



Density-Functional Theory (DFT) + Non-Equilibrium Green's Function (NEGF)

TranSIESTA

J. M. Soler *et al. J. Phys. Condens. Matter.* **14**, 2745 (2002). Mads Brandbyge *et al. Phys. Rev. B* **65**, 165401 (2002). Nick R. Papior. *In preparation* (2016).

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Crossed 14-AGNR



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Crossed 14-AGNR



Crossed 14-AGNR



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Simulation characteristics:

- 1280 atoms;
- double-ζ (9280 orbitals);
- vdW (optB88);
- real space grid cutoff: 350 Ry;
- forces < 5 meV/Å;
- interlayer distance: 3.34 Å.

3.34 Å

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Direct transmission





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Direct transmission





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Direct transmission





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Inter-ribbon transmission



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Electrostatic potential at 0.05 V



 $U_{H}(V=0.05V) (eV)$

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Electrostatic potential at 0.05 V



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Electrostatic potential at 0.05 V



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Electrostatic potential at 0.05 V



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Electrostatic potential at 0.05 V



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Rotated crossbar



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Electrostatic potential at V = 0









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Direct transmission at V = 0





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Direct transmission at V = 0







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Inter-ribbon transmission at V = 0



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Inter-ribbon current



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Inter-ribbon current



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Bond currents for 90° at 0.5 V





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Intralayer bias



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Intralayer bias



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Intralayer bias



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Conclusions

- Application of TranSIESTA for N=4 arbitrarily distributed electrodes at finite bias;
- Transmission **strongly depends on the stacking**;
- For a 60° rotation angle one finds a higher inter-layer transmission;
- In our calculations we observe a **small gating effect** due to the top ribbon.

eman ta zabal zaz





1-D Quantum Well States on Doped Graphene Nanoribbons Revealed by Transport Simulations

Pedro Brandimarte, Aran Garcia-Lekue, Eduard Carbonell-Sanromà, Martina Corso, Richard Balog, Shigeki Kawaii, Shohei Saito, Shinichiro Osumi, Shigehiro Yamaguchi, Jose I. Pascual, and Daniel Sánchez-Portal

June 13, 2016

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Hybrid 7-AGNR

S. Kawai et al. Nature Communications **6**, 8098 (2015). R. R. Cloke et al. J. A. Chem. Soc. **137**, 8872 (2015).

Borinated sectionsPristine sections

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x(nm)

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TranSIESTA



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TranSIESTA



1-D Quantum Well States on Doped Graphene
Nanoribbons Revealed by Transport Simulations

Pedro Brandimarte **TranSIESTA**

Simulation characteristics:

- 756 atoms;
- double-ζ (5040 orbitals);
- vdW (optB88);
- real space grid cutoff: 250 Ry;
- forces < 10 meV/Å.

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Periodic calculation



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Periodic calculation



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Electrostatic potential



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Coulomb cutoff

$$\widetilde{v}(\mathbf{G}) = \int \int \int \int_{space} \widetilde{v}(r) e^{-i\mathbf{G}\cdot\mathbf{r}} d^3\mathbf{r} = \int \int \int \int_{\mathcal{D}} v(r) e^{-i\mathbf{G}\cdot\mathbf{r}} d^3\mathbf{r}$$

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Coulomb cutoff

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• 0-D:
$$\tilde{v}^{0D}(G) = \frac{4\pi}{G^2} [1 - \cos(GR)]$$

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Coulomb cutoff

$$\widetilde{v}(\mathbf{G}) = \int \int \int \int_{space} \widetilde{v}(r) e^{-i\mathbf{G}\cdot\mathbf{r}} d^3\mathbf{r} = \int \int \int \int_{\mathcal{D}} v(r) e^{-i\mathbf{G}\cdot\mathbf{r}} d^3\mathbf{r}$$

• 0-D:
$$\tilde{v}^{0D}(G) = \frac{4\pi}{G^2} [1 - \cos(GR)]$$

• 1-D:
$$\tilde{v}^{1D}(G_x, G_\perp) = \frac{4\pi}{G^2} [1 + G_\perp R J_1(G_\perp R) K_0(G_x R) - G_x R J_0(G_\perp R) K_1(G_x R)]$$

Coulomb cutoff

$$\widetilde{v}(\mathbf{G}) = \int \int \int \int_{space} \widetilde{v}(r) e^{-i\mathbf{G}\cdot\mathbf{r}} d^3\mathbf{r} = \int \int \int \int_{\mathcal{D}} v(r) e^{-i\mathbf{G}\cdot\mathbf{r}} d^3\mathbf{r}$$

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$$\approx -4\pi \int_0^R r J_0(G_\perp r) \ln(r) dr + 4\pi R \ln(2h) \frac{J_1(G_\perp R)}{G_\perp}$$

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Coulomb cutoff

$$\widetilde{v}(\mathbf{G}) = \int \int \int \int_{space} \widetilde{v}(r) e^{-i\mathbf{G}\cdot\mathbf{r}} d^3\mathbf{r} = \int \int \int \int_{\mathcal{D}} v(r) e^{-i\mathbf{G}\cdot\mathbf{r}} d^3\mathbf{r}$$

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$$\approx -4\pi \int_0^R r J_0(G_{\perp}r) \ln(r) dr + 4\pi R \ln(2h) \frac{J_1(G_{\perp}R)}{G_{\perp}}$$

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Electrostatic potential



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DOS projected on each ribbon "row"



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DOS projected on each ribbon "row"



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DOS projected on each ribbon "row"



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Zero bias transmission



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Eigenchannels



M. Paulsson and M. Brandbyge. Phys. Rev. B 76, 115117 (2007).

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Eigenchannels



M. Paulsson and M. Brandbyge. Phys. Rev. B 76, 115117 (2007).

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Conclusions

- Semiconductor electrodes: use with caution!
- **Coulomb cutoff in TranSIESTA** for low dimensionality systems;
- Transport simulations can reproduce observed quantum well states and explain their mechanism.

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Thank you!



